Automatic Calculation of SUSY-particle Production*

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Abstract

We introduce a new method to treat Majorana fermions and interactions with fermionnumber violation on the GRACE system which has been developed for the automatic computation of the matrix elements for the processes of the standard model. Thus we have constructed a system for the automatic computation of cross-sections for the processes of the minimal SUSY standard model (MSSM).

1 Introduction

At the theoretical point of view, it has been a promising hypothesis that there exists a symmetry called supersymmetry (SUSY) between bosons and fermions at the unification-energy scale. It, however, is a broken symmetry at the electroweak-energy scale. Thus the relic of SUSY is expected to remain as a rich spectrum of SUSY particles, partners of usual matter fermions, gauge bosons and Higgs scalars, named sfermions, gauginos and higgsinos, respectively [1].

The quest of these new particles has already been one of the most important pursuits to the present high-energy physics [2]. Although such particles have not yet been discovered, masses of them are expected to be $O(10^2)$ GeV [3]. In order to obtain signatures of the SUSY-particle production, electron-positron colliding experiments are preferable because the electroweak interactions are clean and well-known. Thus we hope SUSY particles will be detected at future e^-e^+ -colliders of sub-TeV-region or TeV-region energies such as LEP2 or NLC's (Next Linear Colliders) [4, 5].

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For the simulations of the experiments, we have to calculate the cross-sections for the processes with the final 3-body or more. We have already known within the standard model that the calculation of the helicity amplitudes is more advantageous to such a case than that of the traces for the gamma matrices with REDUCE [6, 7]. The program package CHANEL [8] is one of the utilities for the numerical calculation of the helicity amplitudes, which has been developing by one of the authors (H.T.).

It, however, is also hard work to construct a program with many subroutine calls of CHANEL by hand. Thus we need a more convenient way to carry out such a work. Several groups have started independently to develop computer systems which automate the perturbative calculation in the standard model with different methods [9, 10, 11, 12, 13]. The GRACE system [9], which automatically generates the source code for CHANEL, is one of the solutions. The system also includes the interface and the library of CHANEL, and the program package BASES/SPRING v5.1 [14] for multi-dimensional integrations and event-generations.

In the SUSY models, there exist Majorana fermions as the neutral gauginos and higgsinos, which become the mixed states called neutralinos. Since anti-particles of Majorana fermions are themselves, there exists so-called 'Majorana-flip', the transition between particle and anti-particle. This is the most important problem which we should solve when we realize the automatic system for computation of the SUSY processes.

In a recent work [15, 16], we developed an algorithm to treat Majorana fermions in CHANEL. In the standard model, we already have such particles as Dirac fermions, gauge bosons and scalar bosons in the GRACE system. There, however, exists another problem on fermion-number violating interactions. We have also developed an algorithm for this problem. Thus we can construct an automatic system for the computation of the SUSY processes by the algorithms above in the GRACE system.

2 SUSY particles and interactions into GRACE 2.0

In Fig. 1, we present the system flow of GRACE (after version 1.1) [17]. The GRACE system has become more flexible for the extension in the new version called 'grc' [18], which is written by C, and includes a new graph-generation package [19]. With this package, any graphs based on a user-defined model can be generated at any orders. The Feynman diagrams are drawn by the program package 'gracefig' [20] in the new GRACE. It is necessary for us to make the interface and the library of CHANEL and the model file for including the SUSY particles.

The method of computation in the program package CHANEL is as follows:

- 1. To divide a helicity amplitude into vertex amplitudes.
- 2. To calculate each vertex amplitude numerically as a complex number.
- 3. To reconstruct of them with the polarization sum, and calculate the helicity amplitudes numerically.

The merit of this method is that the extension of the package is easy, and that each vertex can be defined only by the type of concerned particles.

Here we use an algorithm [15, 16] for the implementation of the embedding Majorana fermions in CHANEL as follows:

• policy

- 1. To calculate a helicity amplitude numerically.
- 2. To replace each propagator by wave functions or polarization vectors, and calculate vertex amplitudes.
- 3. **Not to** move charge-conjugation matrices.

• method

- 1. To choose a direction on a fermion line.
- 2. To put wave functions, vertices and propagators along the direction in such a way:
 - i) To take the transpose for the reverse direction of fermions

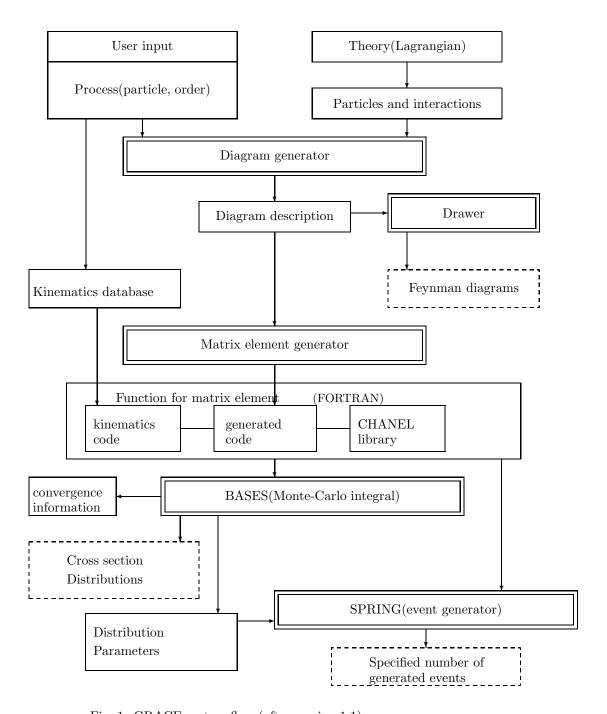


Fig. 1. GRACE system flow (after version 1.1)

ii) To use the propagator with the charge-conjugation matrix for the Majorana-flipped one.

As a result, the kinds of the Dirac-Majorana-scalar vertices are limited to four types:

$$J_{1\ h_1 h_2}^{[S]\rho_1 \rho_2} = \overline{U}^{\rho_1}(h_1, p_1, m_1) \Gamma U^{\rho_2}(h_2, p_2, m_2) , \qquad (2.1)$$

$$J_{1\ h_{1}h_{2}}^{[S]\rho_{1}\rho_{2}} = \overline{U}^{\rho_{1}}(h_{1}, p_{1}, m_{1})\Gamma U^{\rho_{2}}(h_{2}, p_{2}, m_{2}) ,$$

$$J_{2\ h_{1}h_{2}}^{[S]\rho_{1}\rho_{2}} = U^{\rho_{1}T}(h_{1}, p_{1}, m_{1})\Gamma \overline{U}^{\rho_{2}T}(h_{2}, p_{2}, m_{2}) ,$$

$$(2.1)$$

$$J_{3 h_1 h_2}^{[S]\rho_1 \rho_2} = \overline{U}^{\rho_1}(h_1, p_1, m_1)C^{\mathrm{T}}\Gamma^{\mathrm{T}}\overline{U}^{\rho_2 \mathrm{T}}(h_2, p_2, m_2) ,$$

$$J_{4 h_1 h_2}^{[S]\rho_1 \rho_2} = U^{\rho_1 \mathrm{T}}(h_1, p_1, m_1)\Gamma^{\mathrm{T}}C^{-1}U^{\rho_2}(h_2, p_2, m_2) ,$$

$$(2.3)$$

$$J_{4\ h_1 h_2}^{[S]\rho_1 \rho_2} = U^{\rho_1 T}(h_1, p_1, m_1) \Gamma^T C^{-1} U^{\rho_2}(h_2, p_2, m_2) , \qquad (2.4)$$

where U's denote wave functions, and C is the charge-conjugation matrix. The symbol Γ stands for the scalar vertex such as

$$\Gamma = A_{\rm L} \cdot \frac{1 - \gamma_5}{2} + A_{\rm R} \cdot \frac{1 + \gamma_5}{2} \quad .$$

The vertices $J_2^{[{\rm S}]} \sim J_4^{[{\rm S}]}$ are related to the vertex $J_1^{[{\rm S}]}$ which has been already defined as the Dirac-Dirac-scalar vertex in the subroutine FFS of CHANEL. The relations among the vertices are as follows:

$$J_{1\ h_1 h_2}^{[S]\rho_1\rho_2} \longrightarrow FFS , \qquad (2.5)$$

$$J_{1 h_{1} h_{2}}^{S[]} \longrightarrow FFS , \qquad (2.5)$$

$$J_{2 h_{1} h_{2}}^{[S] \rho_{1} \rho_{2}} = -J_{1 h_{1} h_{2}}^{[S] - \rho_{1} - \rho_{2}} \longrightarrow FFST , \qquad (2.6)$$

$$J_{3 h_{1} h_{2}}^{[S] \rho_{1} \rho_{2}} = -J_{1 h_{1} h_{2}}^{[S] \rho_{1} - \rho_{2}} \longrightarrow FFCS , \qquad (2.7)$$

$$J_{4 h_{1} h_{2}}^{[S] \rho_{1} \rho_{2}} = -J_{1 h_{1} h_{2}}^{[S] - \rho_{1} \rho_{2}} \longrightarrow FFSC , \qquad (2.8)$$

$$J_{3\ h_1h_2}^{[S]\rho_1\rho_2} = -J_{1\ h_1h_2}^{[S]\rho_1-\rho_2} \longrightarrow FFCS ,$$
 (2.7)

$$J_{4\ h_1 h_2}^{[S]\rho_1\rho_2} = -J_{1\ h_1 h_2}^{[S]-\rho_1\rho_2} \longrightarrow FFSC ,$$
 (2.8)

Thus we can build three new subroutines FFST, FFCS and FFSC. We have performed the installation of the subroutines above with the interface on the GRACE system version 2.0 [21, 16, 22, 23, 24, 25].

Next we propose an algorithm for the interactions with fermion-number violation such as the charginoselectron-antineutrino vertex. We introduce two new soubroutines [26, 27].

$\mathbf{3}$ Tests for the system

At the first stage for the check of our system, we have written the model file of the SUSY QED. In this case, there is only one Majorana fermion called photino. It is essential for testing our system to include photino and its interactions.

Next we have extended the model file and the definition file of couplings for the MSSM. The tests have been performed by the exact calculations with the two methods, our system and REDUCE, in such a manner:

- 1. To calculate the differential cross-sections at a point of the phase space in the two methods with GRACE and REDUCE.
- 2. To calculate the differential cross-sections over the phase space in the two methods with GRACE and REDUCE through BASES.
- 3. To integrate the differential cross-sections over the phase space in the two methods with GRACE and REDUCE through BASES.

With BASES, we can get the differential cross-sections and the scattered plots by one time of the integration step. For writing REDUCE sources, we use the different method to treat Majorana fermions in Ref. [13].

In Table I, the tested processes are shown as a list. The references in the table (without [16], [24] and [32]) are not the results of the tests, but for help.

In Ref. [32], we show the angular distribution of the outgoing positron in the process $e^-e^+ \to \tilde{e}_B^- \tilde{\gamma} e^+$. Here we use BASES for the calculation from the REDUCE output. The result is in beautiful agreement with the value that is obtained by GRACE at each bin of the histogram. Since the two diagrams with the one-photon exchange dominate in this case, there is a steep peak in the direction of the initial positron. In such a case, the equivalent-photon approximation (EPA) works well [31].

Process		Number of diagrams	Comment	Check	Reference
SUSY	QED				
$e^-e^- \rightarrow$	$ \tilde{e}_{\mathrm{R}}^{-}\tilde{e}_{\mathrm{R}}^{-} $ $ \tilde{e}_{\mathrm{L}}^{-}\tilde{e}_{\mathrm{L}}^{-} $	2	Majorana-flip	OK	[16]
	$ ilde{e}_{ m L}^{-} ilde{e}_{ m L}^{-}$	2	in internal lines	OK	[16]
	$\tilde{e}_{\mathrm{R}}^{-}\tilde{e}_{\mathrm{L}}^{-}$	2		OK	[16]
$e^-e^+ \rightarrow$	$\tilde{e}_{\mathrm{R}}^{-}\tilde{e}_{\mathrm{R}}^{+}$	2	Including pair	OK	[23, 28]
	$ ilde{e}_{ m L}^{-} ilde{e}_{ m L}^{+}$	2	annihilation	OK	[23, 28]
$e^-e^+ \rightarrow$	$ ilde{e}_{ m R}^- ilde{e}_{ m L}^+ \ ilde{e}_{ m R}^+ ilde{e}_{ m L}^-$	1	Values are	OK	[23, 28]
	$ ilde{e}_{ m R}^{\mp} ilde{e}_{ m L}^{\mp}$	1	equal	OK	[23, 28]
$e^-e^+ \rightarrow$	$\tilde{\gamma}\tilde{\gamma}$	4	F-B symmetric	OK	[16]
$e^-e^+ \rightarrow$	$\tilde{\gamma}\tilde{\gamma}\gamma$	12	Final 3-body	OK	[29]
$e^-e^+ \rightarrow$	$\tilde{e}_{\mathrm{R}}^{-}\tilde{\gamma}e^{+}$	12	Including every	OK	[30, 31]
			elements for tests		[32]
MSSM		_	-		
$e^-e^- \rightarrow$	$\tilde{e}_{\mathrm{L}}^{-}\tilde{e}_{\mathrm{L}}^{-}$	8	4 Majorana fermions	OK	[24]
$e^-e^+ \rightarrow$	$\tilde{\chi_1}^-\tilde{\chi_1}^+$	3		OK	[24]
$e^-e^+ \rightarrow$	$\tilde{\chi}_1^0 \tilde{\chi}_1^0$	14	Final 3-body	OK	[24]

Table I. The list of the tested processes.

4 Summary

We introduce a new method to treat Majorana fermions and interactions with fermion-number violation on the GRACE system for the automatic computation of the matrix elements for the processes of the SUSY models. In the first instance, we have constructed the system for the processes of the SUSY QED because we should test our algorithm for the simplest case. Next we have extended the model file and the definition file of couplings for the MSSM. The numerical results convince us that our algorithm is correct.

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